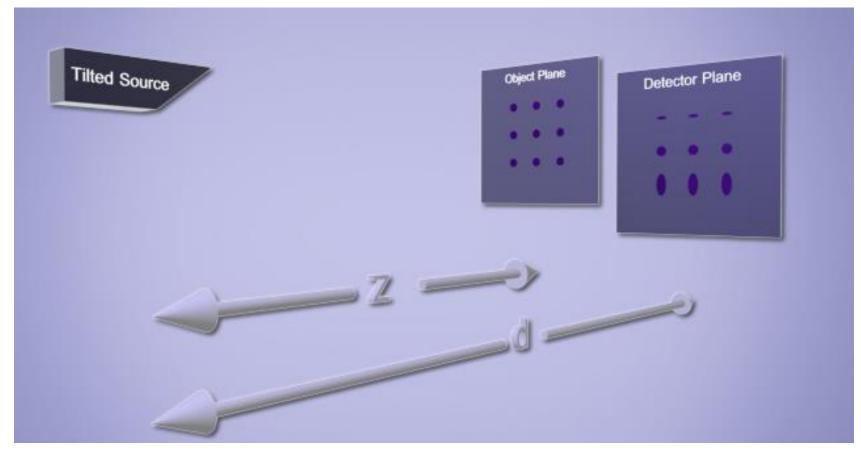
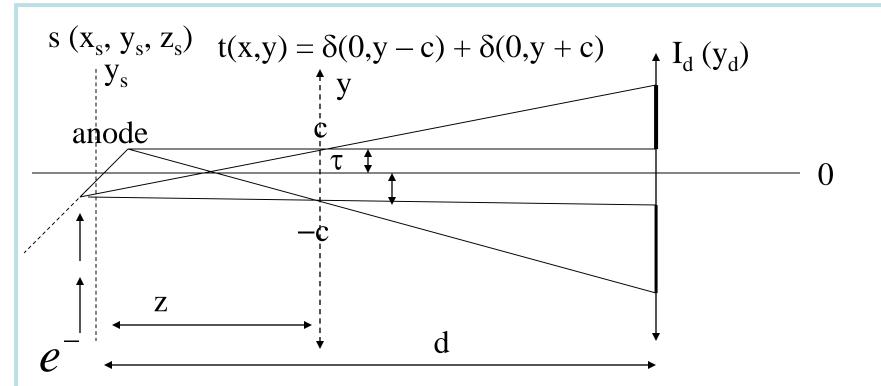
Effects of Tilted Source (Anode Angle)

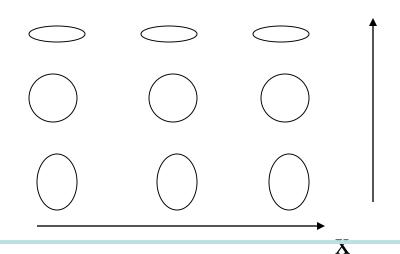


Source magnification in x_d function of z Source magnification in y dependent on y_d and z

We will skip the mathematical development on this section

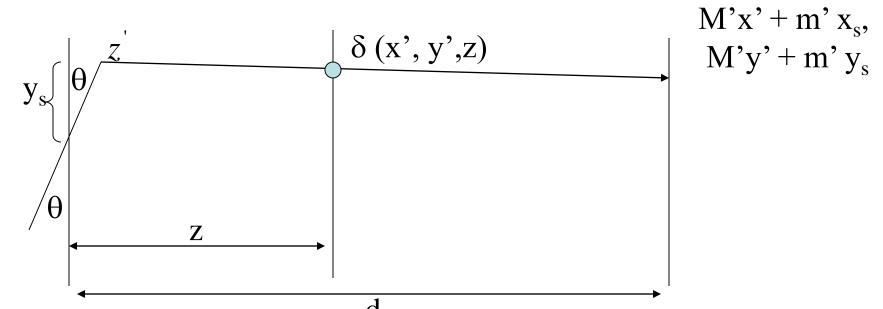


Consider a 3 x 3 array of pinholes



t(x,y) above can be considereda lead plate with two pinholes punchedv into it.

3 x 3 array of circular pinholes to left shows how source is contracted in y for positive y detector positions and enlarged for negative detector locations. Let's allow the magnification to be different on each axis. The pinhole models an object, $\delta(x', y', z)$.



The diagram above merely shows how one point in the source the pinhole, and the detector are related by geometry.

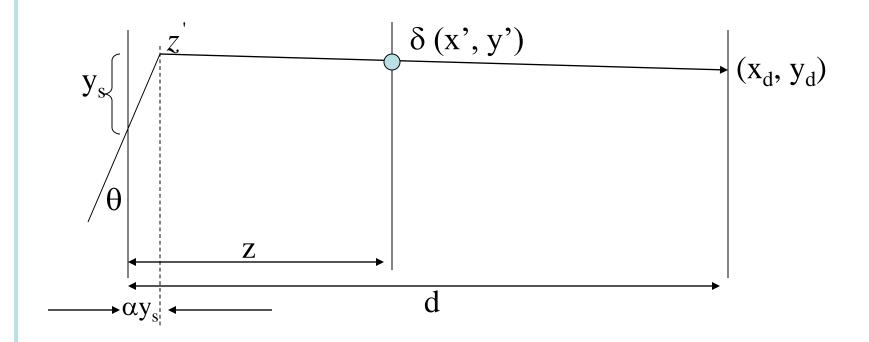
Here we use a M' to allow for magnification of the pinhole, and m' to allow for magnification of the source

We will consider the impulse response using a pinhole at x', y'.

$$h(x_{d}, y_{d}, x', y') = \frac{\eta}{m_{x}m_{y}} s(\frac{x_{d} - M_{x}x'}{m_{x}}, \frac{y_{d} - M_{y}y'}{m_{y}})$$

where η is a collection efficiency for the pinhole: $\eta = \Omega/4\pi d^2$

If $tan (\theta) = \alpha$, then the source position $z' = \alpha y_s$



By geometry, object magnification is

$$M' = (d - \alpha y_s)/(z - \alpha y_s)$$

Source magnification is

$$m' = -((d - \alpha y_s) - (z - \alpha y_s)) / (z - \alpha y_s)$$
$$m' = -(d - z) / (z - \alpha y_s)$$

Recall $x_d = M'x' + m'x_s$ $y_{d} = M'y' + m'y_{s}$ $M_x = \partial x_d / \partial x' = M' = (d - \alpha y_s) / (z - \alpha y_s) \approx d/z$ since for practical arrangements d, $z >> \alpha y_s$ Typical dimensions: z, d ~ 1 m, y_s ~ 1mm Similarly $M_v = \partial y_d / \partial y' \approx d/z$ $m_x = \partial x_d / \partial x_s = m' \approx - (d-z)/z = m$ $m_v = \partial y_d / \partial y_s$ This is more interesting derivative since both M' and m' are functions of y_s

 $\mathbf{m}_{\mathbf{y}} = \partial \mathbf{y}_{\mathbf{d}} / \partial \mathbf{y}_{\mathbf{s}} = \partial (\mathbf{M}' \mathbf{y}') / \partial \mathbf{y}_{\mathbf{s}} + \partial (\mathbf{m}' \mathbf{y}_{\mathbf{s}}) / \partial \mathbf{y}_{\mathbf{s}}$

From previous slide,

M' =
$$(d - \alpha y_s)/(z - \alpha y_s)$$
 and **m**' = $-(d - z)/(z - \alpha y_s)$

To find:
$$m_y = \partial y_d / \partial y_s = \partial (M' y') / \partial y_s + \partial (m' y_s) / \partial y_s$$

$$m_{y} = [((z - \alpha y_{s})(-\alpha) - (d - \alpha y_{s})(-\alpha)] y')/(z - \alpha y_{s})^{2} + -(d-z) \cdot [(z - \alpha y_{s}) - y_{s}(-\alpha)]/(z - \alpha y_{s})^{2} = (-\alpha[z-d] y' - (d - z) z)/(z - \alpha y_{s})^{2} = -(d - z) (z - \alpha y')/((z - \alpha y_{s})^{2})^{2}$$

$$m_y = - (d-z) (z - \alpha y') / z^2 = m (1 - (\alpha y'/z))$$

Using this relationship and ignoring obliquity,

$$h(x_d, y_d, x', y') = \frac{1}{4\pi d^2 m^2} \frac{1}{(1 - \alpha y'/z)} s(\frac{x_d - M x'}{m}, \frac{y_d - M y'}{m(1 - \alpha y'/z)})$$

How does magnification change with object position?

Since system is linear, we can write a superposition integral.

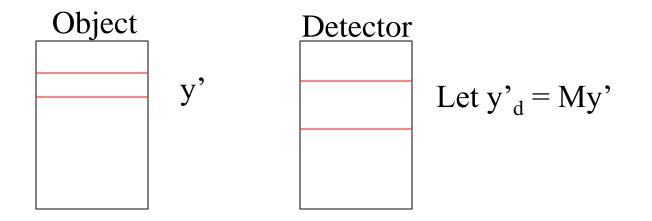
$$I_{d}(x_{d}, y_{d}) = \int \int h((x_{d}, y_{d}, x', y') t(x', y') dx' dy' =$$
$$\iint \frac{1}{4\pi d^{2}m^{2}} \frac{1}{(1 - \alpha y'/z)} s(\frac{x_{d} - M x'}{m}, \frac{y_{d} - M y'}{m(1 - \alpha y'/z)}) t(x', y') dx' dy'$$

For developing space-invariance, let's consider a magnified object

Let
$$x'' = Mx'$$
 $y'' = My'$
 $dx'' = Mdx'$ $dy'' = Mdy'$
 $I_d(x_d, y_d) = \frac{1}{4\pi d^2 m^2} \iint \frac{1}{(1 - \alpha y'' / Mz)} s(\frac{x_d - x''}{m}, \frac{y_d - y''}{m(1 - \alpha y'' / Mz)})t(\frac{x''}{M}, \frac{y''}{M})dx''dy''$

Not a space-invariant system since $1 - (\alpha y''/Mz)$ varies slowly with y'' or y'

But it doesn't vary much in a region of an object.

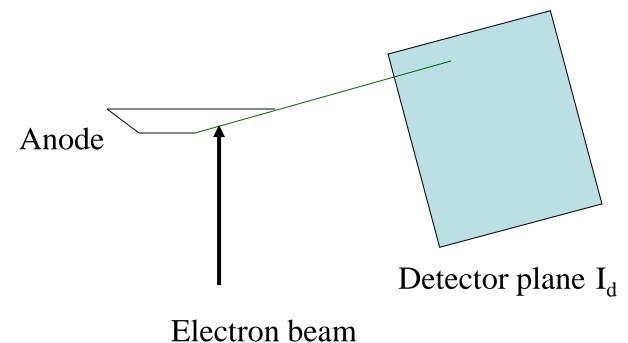


Consider a horizontal strip across the detector centered at For this region, $\alpha y'' = \alpha M y' = \alpha y_d$ where Mz = d

$$I_{d}(x_{d}, y_{d}, y_{d'}) = \frac{1}{4\pi d^{2}m^{2}(1 - \alpha y_{d'}/d)} s(\frac{x_{d}}{m}, \frac{y_{d}}{m(1 - \alpha y_{d'}/d)}) * t(\frac{x_{d}}{M}, \frac{y_{d}}{M})$$

Here y'_d is a constant over a region in the detector during the convolution.

At $\alpha y' = z$, source width goes to 0 in y. We call this the "heel effect"

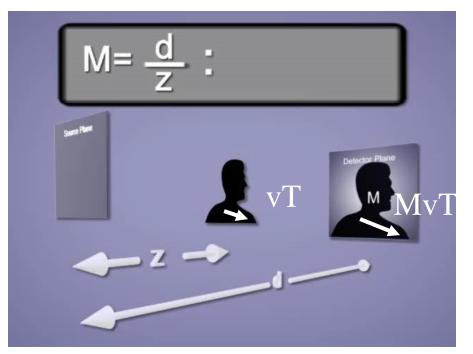


In the frequency domain,

$$I_{d}(u,v,y_{d'}) = \frac{M^{2}}{4\pi d^{2}}S(mu,m(1-\frac{\alpha y_{d}}{d})v)T(Mu,Mv)$$



Let's consider object motion with constant velocity in the x direction over the imaging time T at velocity v



Over time T, the object size will change position in the detector plane by MvT

The impulse response due to movement in x is

$$h_{movement}(x_d, y_d) = \frac{1}{MvT} \prod(\frac{x_d}{MvT}) \delta(y_d)$$

Notice that there is no degradation in y, as we expected.

The complete impulse response is given below(a planar source parallel to the detector is used here for simplicity).

$$I_{d} = \frac{1}{4\pi d^{2}m^{2}}t(\frac{x_{d}}{M}, \frac{y_{d}}{M}) * s(\frac{x_{d}}{m}, \frac{y_{d}}{m}) * \frac{1}{MvT}\Pi(\frac{x_{d}}{MvT})$$

Blur in x direction gets minimized as T decreases to 0

How to minimize Motion blurring

Assume a L x L source parallel to detector

$$s(x_s, y_s) = K \prod(\frac{x_s}{L}) \prod(\frac{y_s}{L})$$

Write energy density E_s as a power integrated over time

 $p(x_s, y_s)$ is regional source power density (it is limited by tungsten melting point). Set $p(x_s, y_s) = P_{max.}$ Operating tube at maximum power available. T is the time beam is ON

$$E_{s} = \iiint p(x_{s}, y_{s}) dx_{s} dy_{s} dt = P_{max}TL^{2} , \text{ then } T = \frac{E_{s}}{P_{max}L^{2}}$$

If L increases, source grows, then T can decreases.

Complete Response functionhfor Rectangle Source andmovement MvT in x direction:

$$h(x_d, y_d) = K \prod(\frac{x_d}{mL}, \frac{y_d}{mL}) * * \prod(\frac{x_d}{MvE_s})$$

$$\frac{MvE_s}{D}$$

I max L

We have extension of the impulse response in the x direction as: $X = |m|L + (MvE_s/P_{max} L^2)$

by the convolution of two rectangles, due to source blurring and motion

We could choose to minimize several criteria. Area, for example. We will simply minimize X now with respect to L.

$$L_{\min} = \left| \frac{2MvE_s}{P_{\max}|m|} \right|^{1/3}$$

Corresponding Exposure Time at Optimal L

$$T = \frac{E_s}{P_{\text{max}}L^2} = \frac{E_s}{P_{\text{max}}} \left[\frac{P_{\text{max}} |m|}{2MvE_s} \right]^{2/3} = \left[\frac{E_s}{P_{\text{max}}} \right]^{1/3} \left[\frac{|m|}{2Mv} \right]^{2/3}$$

If v = 0, thenL = 0 $T = \infty$ no source blurringIf |m| = 0,object is on detector $L \rightarrow \infty$ $\rightarrow T = 0$